

***D*-Dimensional *q*-Harmonic Oscillator and *d*-Dimension *q*-Hydrogen Atom**

Su Ka-Lin¹ and Liu An-ling²

Received May 1, 1999

A *q*-analogue of the *D*-dimensional harmonic oscillator is presented. A new realization of the quantum algebra $SU_q(1,1)$ via the *D*-dimensional *q*-harmonic oscillator is found. A model of the *d*-dimensional *q*-hydrogen atom is constructed by means of the *D*-dimensional *q*-harmonic oscillator. The dimension *D* of the *q*-harmonic oscillator and the dimension *d* of the *q*-hydrogen atom are arbitrary.

In recent years, the *q*-analogue of the one-dimensional harmonic oscillator has been studied by several authors [1–3]. Realizations of the quantum algebra $SU_q(1,1)$ via the one-dimensional *q*-harmonic oscillator were suggested by Chaichian *et al.* [4, 5]. A model of the three-dimensional *q*-hydrogen atom has been constructed by Kibler and Negadi [6] and others [7, 8]. In this paper, we will present the *q*-analogue of the *D*-dimensional harmonic oscillator, find a new realization of the quantum algebra $SU_q(1,1)$ via the *D*-dimensional *q*-harmonic oscillator, and construct a model of the *D*-dimensional *q*-hydrogen atom by means of the *D*-dimensional *q*-harmonic oscillator. The dimension *D* of the *q*-harmonic oscillator and the dimension *d* of the *q*-hydrogen atom are arbitrary.

1. THE *q*-ANALOGUE OF THE *D*-DIMENSIONAL HARMONIC OSCILLATOR

The annihilation and creation operators of the *D*-dimensional *q*-harmonic oscillator $a_{q\alpha}$ and $a_{q\alpha}^+$ ($\alpha = 1, 2, 3, \dots, D$) can be defined as

¹Department of Physics, Yueyang Normal College, Yueyang 414000, China.

²Department of Physics, Changsha Electric Power College, Changsha 410000, China.

$$a_{q\alpha} = \left\{ \frac{[N_\alpha + 1]}{N_\alpha + 1} \right\}^{1/2} a_\alpha, \quad a_{q\alpha}^+ = a_\alpha^+ \left\{ \frac{[N_\alpha + 1]}{N_\alpha + 1} \right\}^{1/2} \tag{1}$$

where a_α and a_α^+ are the annihilation and creation operators of the D -dimensional harmonic oscillator, $N_\alpha = a_\alpha^+ a_\alpha$, and $[x] = (q^x - q^{-x})/(q - q^{-1})$. These operators satisfy the following relations:

$$a_{q\alpha}^+ a_{q\alpha} = [N_\alpha], \quad a_{q\alpha} a_{q\alpha}^+ = [N_\alpha + 1], \quad a_{q\alpha} a_{q\alpha}^+ - q a_{q\alpha}^+ a_{q\alpha} = q^{-N_\alpha} \tag{2}$$

$$[a_{q\alpha}, a_{q\beta}] = [a_{q\alpha}^+, a_{q\beta}^+] = [N_\alpha, N_\beta] = 0 \tag{3}$$

where $\alpha = 1, 2, 3, \dots, D$ and $\beta = 1, 2, 3, \dots, D$. The Hamiltonian H'_q of the D -dimensional q -harmonic oscillator can be defined as

$$H'_q = \frac{1}{2} \sum_{\alpha=1}^D (a_{q\alpha}^+ a_{q\alpha} + a_{q\alpha} a_{q\alpha}^+) = \frac{1}{2} \sum_{\alpha=1}^D ([N_\alpha] + [N_\alpha + 1]) \tag{4}$$

where we have assumed $\hbar\omega = 1$; the eigenequation of H'_q is

$$H'_q |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle = E'_{qN} |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle \tag{5}$$

where E'_{qN} is the energy eigenvalue of H'_q , and $|n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle$ is the corresponding eigenvector. Because

$$N_\alpha |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle = n_\alpha |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle \tag{6}$$

we have

$$[N_\alpha] |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle = [n_\alpha] |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle \tag{7}$$

Equation (5) becomes

$$\begin{aligned} & H'_q |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle \\ &= \left\{ \frac{1}{2} \sum_{\alpha=1}^D ([n_\alpha] + [n_\alpha + 1]) \right\} |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle \end{aligned} \tag{8}$$

Then we have

$$E'_{qN} = \frac{1}{2} \sum_{\alpha=1}^D ([n_\alpha] + [n_\alpha + 1]) \tag{9}$$

Equation (9) is the energy spectrum formula of the D -dimensional q -harmonic oscillator.

2. THE *D*-DIMENSIONAL *q*-HARMONIC OSCILLATOR REALIZATION OF THE QUANTUM ALGEBRA $SU_q(1,1)$

We use the annihilation and creation operators $a_{q\alpha}$ and $a_{q\alpha}^+$ ($\alpha = 1, 2, 3, \dots, D$) of the *D*-dimensional *q*-harmonic oscillator to construct the operators K'_1, K'_2 , and K'_3 .

$$\begin{aligned} K'_1 &= \frac{1}{4} \sum_{\alpha=1}^D [(a_{q\alpha}^+)_2 + a_{q\alpha}^2] \\ K'_2 &= -\frac{i}{4} \sum_{\alpha=1}^D [(a_{q\alpha}^+)_2 - a_{q\alpha}^2] \\ K'_3 &= \frac{1}{4} \sum_{\alpha=1}^D (a_{q\alpha}^+ a_{q\alpha} + a_{q\alpha} a_{q\alpha}^+) \end{aligned} \tag{10}$$

Then we define K'_+ and K'_- by $K'_\pm = K'_1 \pm iK'_2$; it is easy to show that the operators K'_+, K'_- , and K'_3 satisfy the following commutation relations:

$$[K'_3, K'_\pm] = \pm K'_\pm, \quad [K'_+, K'_-] = -[2K'_3] \tag{11}$$

These are the familiar commutation relations of the quantum algebra $SU_q(1,1)$. This indicates that we have realized the quantum algebra $SU_q(1,1)$ via the *D*-dimensional *q*-harmonic oscillator.

3. MODEL OF THE *d*-DIMENSIONAL *q*-HYDROGEN ATOM

In order to construct a model of the *d*-dimensional *q*-hydrogen atom, first we find the relationship between the *d*-dimensional hydrogen atom and the *D*-dimensional harmonic oscillator [9, 10]. We use $x_1, x_2, x_3, \dots, x_d$ to express the *d*-dimensional hydrogen atom coordinate space and construct the operators K_1, K_2 , and K_3 for the space [11]:

$$K_1 = \frac{1}{2} (r\Delta + r), \quad K_2 = i \left[\frac{d-1}{2} + \sum_{j=1}^d x_j \frac{\partial}{\partial x_j} \right], \quad K_3 = -\frac{1}{2} (r\Delta - r) \tag{12}$$

where $\Delta = \sum_{j=1}^d \partial^2 / \partial x_j \partial x_j$ and $r = (\sum_{j=1}^d x_j^2)^{1/2}$. It is easy to show that the operators satisfy the following commutation relations:

$$[K_1, K_2] = -iK_3, \quad [K_2, K_3] = iK_1, \quad [K_3, K_1] = iK_2 \tag{13}$$

These relations show that the operators K_1, K_2 , and K_3 constitute the $SU(1,1)$ algebra. We write the Hamiltonian of the *d*-dimensional hydrogen atom as

$$H = -\frac{1}{2} \Delta - \frac{1}{r} \quad (14)$$

where we have assumed $\hbar = \mu = e = 1$. Using Eq. (12), we can reduce Eq. (14) to

$$(K_1 + K_3)H = -\frac{1}{2} (K_1 - K_3) - 1 \quad (15)$$

The eigenequation of the Hamiltonian H and eigenvalue can be written as [8, 9, 12, 13]

$$H|d, n\rangle = E_n|d, n\rangle \quad (16)$$

$$E_n = -\frac{1}{2} \frac{1}{[n + (d-3)/2]^2} \quad (17)$$

where E_n is the energy eigenvalue of the Hamiltonian H , and $|d, n\rangle$ is the corresponding eigenvector. From Eqs. (15) and (16), we obtain

$$\left\{ -\left(\frac{1}{2} + E_n\right)K_1 + \left(\frac{1}{2} - E_n\right)K_3 - 1 \right\} |d, n\rangle = 0 \quad (18)$$

Defining the function θ_n by

$$\cosh \theta_n = \frac{1 - 2E_n}{\sqrt{-8E_n}}, \quad \sinh \theta_n = -\frac{1 + 2E_n}{\sqrt{-8E_n}} \quad (19)$$

and using the relation satisfied by the elements of the $SU(1,1)$ algebra

$$e^{-iK_2\theta_n} K_3 e^{iK_2\theta_n} = K_3 \cosh \theta_n + K_1 \sinh \theta_n \quad (20)$$

we can be rewrite Eq. (18), as

$$\left\{ K_3 - \frac{1}{\sqrt{-2E_n}} \right\} e^{-iK_2\theta_n} |d, n\rangle = 0 \quad (21)$$

This is an eigenequation of the operator K_3 . Thus, we have transformed the eigenequation of the Hamiltonian of the d -dimensional hydrogen atom into an eigenequation of the operator K_3 .

We use the annihilation and creation operators a_α and a_α^+ ($\alpha = 1, 2, 3, \dots, D$) of the D -dimensional harmonic oscillator to realize the $SU(1,1)$ algebra given in ref. 14. The results are

$$\begin{aligned} K_1 &= \frac{1}{4} \sum_{\alpha=1}^D [(a_\alpha^+)^2 + a_\alpha^2] \\ K_2 &= -\frac{i}{4} \sum_{\alpha=1}^D [(a_\alpha^+)^2 - a_\alpha^2] \end{aligned} \quad (22)$$

$$K_3 = \frac{1}{4} \sum_{\alpha=1}^D (a_{\alpha}^+ a_{\alpha} + a_{\alpha} a_{\alpha}^+)$$

obviously K_1 , K_2 , and K_3 satisfy Eq. (13). In other words, K_1 , K_2 , and K_3 constitute the $SU(1,1)$ algebra. The Hamiltonian H' of the D -dimensional harmonic oscillator is

$$H' = \frac{1}{2} \sum_{\alpha=1}^D (a_{\alpha}^+ a_{\alpha} + a_{\alpha} a_{\alpha}^+) \tag{23}$$

where we have assumed $\omega\hbar = 1$. The eigenequation and the eigenvalue of the operator H' can be written as

$$H'|n_1, n_2, \dots, n_{\alpha}, \dots, n_D\rangle = E'_N|n_1, n_2, \dots, n_{\alpha}, \dots, n_D\rangle \tag{24}$$

$$E'_N = N + D/2 \tag{25}$$

where E'_N is the energy eigenvalue of the operator H' , $|n_1, n_2, \dots, n_{\alpha}, \dots, n_D\rangle$ is the corresponding eigenvector, and N is the eigenvalue of the operator $\sum_{\alpha=1}^D a_{\alpha}^+$, $N = n_1 + n_2 + \dots + n_D$. Comparing Eq. (22) with Eq. (23), we can rewrite Eq. (24) as

$$K_3|n_1, n_2, \dots, n_{\alpha}, \dots, n_D\rangle = \frac{1}{2} E'_N|n_1, n_2, \dots, n_{\alpha}, \dots, n_D\rangle \tag{26}$$

This also is an eigenequation of the operator K_3 . Comparing Eq. (21) with Eq. (26), one can find relations

$$|d, n\rangle = e^{-iK_2\theta_n}|n_1, n_2, \dots, n_{\alpha}, \dots, n_D\rangle \tag{27}$$

$$E_N = -\frac{2}{(E'_N)^2} \tag{28}$$

$$D = \frac{D}{2} + 1, \quad n = \frac{n}{2} + 1 \tag{29}$$

where

$$K_2 = -\frac{i}{4} \sum_{\alpha=1}^D [(a_{\alpha}^+)^2 - a_{\alpha}^2]$$

Equations (27)–(29) relate the d -dimensional hydrogen atom and the D -dimensional harmonic oscillator. From Eqs. (27)–(29) and (9), we find the energy spectrum of the d -dimensional q -hydrogen atom

$$E_{q_n, n_1 n_2 \dots n_D} = -\frac{2}{(E'_{qN})^2} = -\frac{8}{\{\sum_{\alpha=1}^D ([n_\alpha] + [n_\alpha + 1])\}^2} \quad (30)$$

where $D = 2(d - 1)$, $N = 2(n - 1) = n_1 + n_2 + \dots + n_D$, and the corresponding eigenvector is

$$|d, n\rangle_{q, n_1 n_2 \dots n_D} = e^{-iK'_2 \theta_{qn}} |n_1, n_2, \dots, n_\alpha, \dots, n_D\rangle \quad (31)$$

where

$$K'_2 = -\frac{i}{4} \sum_{\alpha=1}^D [(a_{q\alpha}^+)^2 - a_{q\alpha}^2]$$

The function θ_{qn} is defined by

$$\cosh \theta_{qn} = \frac{1 - 2E_{qn}}{\sqrt{-8E_{qn}}}, \quad \sinh \theta_{qn} = -\frac{1 - 2E_{qn}}{\sqrt{-8E_{qn}}} \quad (32)$$

The d -dimensional q -hydrogen atom energy spectrum (30) shows that the d -dimensional q -hydrogen atom has the same ground energy level as the ordinary d -dimensional hydrogen atom; the excited energy levels ($n > 2$) of the d -dimensional q -hydrogen atom are split. For example, when $n = 2$, $d = 3$, we have

$$E_{q2,2000} = -\frac{8}{([2] + [3] + 3)^2}, \quad E_{q2,1100} = -\frac{2}{([2] + 2)^2} \quad (33)$$

and when $n = 3$, $d = 3$, we have

$$\begin{aligned} E_{q3,1111} &= -\frac{1}{2([2] + 1)^2}, & E_{q3,2200} &= -\frac{2}{([2] + [3] + 1)^2} \\ E_{q3,3100} &= -\frac{8}{([3] + [4] + [2] + 3)^2}, & E_{q3,2110} &= -\frac{8}{(3[2] + [3] + 3)^2} \\ E_{q3,4000} &= -\frac{1}{([4] + [5] + 3)^2} \end{aligned} \quad (34)$$

Equations (30) and (31) show that we have constructed a model of the d -dimensional q -hydrogen atom; when $d = 3$, Eq. (30) is just Gora's result [7] when $q \rightarrow 1$, the results (30) and (31) become the classical case [see (27) and (28)]. Therefore, our work is consistent.

REFERENCES

1. Le-Man Kuang, *J. Phys. A* **26**, L1079 (1993).
2. A. S. Zhedanov, *J. Phys. A* **25**, L713 (1992).

3. A. J. Macfarlane, *J. Phys. A* **22**, 4581 (1989).
4. M. Chaichian *et al.*, *Phys. Rev. Lett.* **65**, 980 (1990).
5. J. Katriel and A. I. Solomon, *J. Phys. A* **24**, 2093 (1991).
6. M. Kibler and T. Negadi, *J. Phys. A* **24**, 5283 (1991).
7. J. Gora, *J. Phys. A* **25**, L1281 (1992).
8. Song Xing-Chang and Liao Li, *J. Phys. A* **25**, 623 (1992).
9. Zeng Gao-Jian, Su Ka-lin, and Li Min, *Phys. Rev. A* **50**, 4373 (1994).
10. Zeng Gao-Jian, Su Ka-Lin, and Li Min, *Chin. Phys. Lett.* **11**, 724 (1994).
11. Su Ka-Lin, *Chin. Phys. Lett.* **14**, 721 (1997).
12. A. O. Barut and I. H. Duru, *Proc. R. Soc. A* **333**, 217 (1973).
13. M. M. Nieto, *Am. J. Phys.* **47**, 1067 (1979).
14. T. Lisowski, *J. Phys. A* **25**, 1259, (1992).